
Enhanced Resistance during Neutron Emission in Wire-Array Z-Pinch

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Outline



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 - Plasma resistance**
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- Summary and conclusion

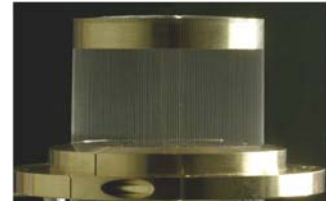
Introduction

Purpose of our experiment



Wire-array Z-pinch

- **Powerful sources of soft X-rays**
(290 TW, 2 MJ, ICF driver)
⇒ detailed studies of EUV, soft and hard X-rays
- **Information about fast ions is rather rare**



<http://zpinch.sandia.gov/>

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Ion Viscous Heating in a Magneto-hydrodynamically Unstable Z Pinch at Over 2×10^9 Kelvin

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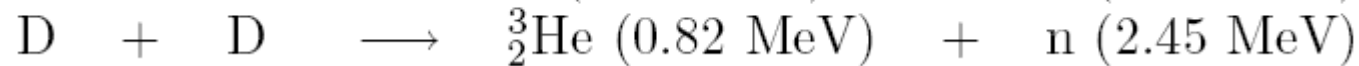
Introduction

Purpose of our experiment



Our primary interest

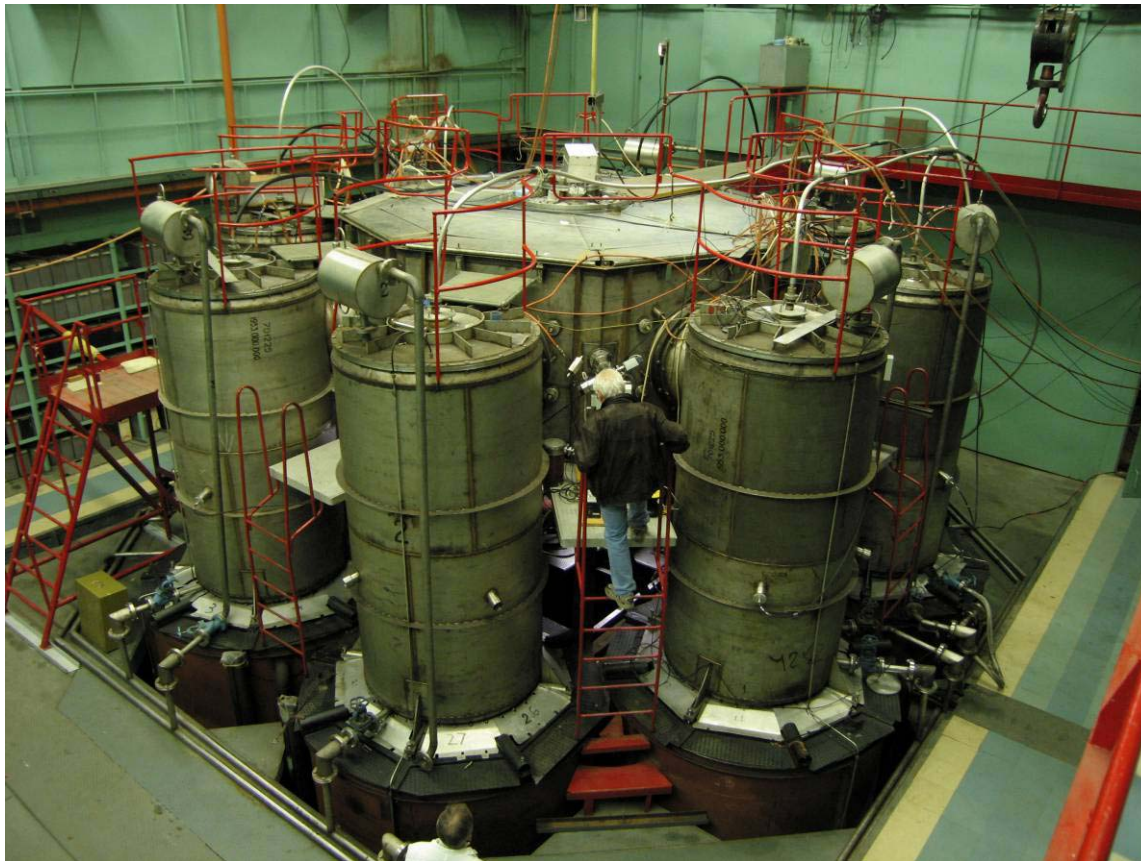
- **To perform fusion neutron measurement in Z-pinches**
⇒ to put a deuterated fibre in the centre of a wire array.



- **Neutrons** ⇒ **favourable diagnostic tool of D-D reaction**
 - = are influenced neither by E nor B
 - = weak absorption and scattering in the chamber walls and air (TOF analysis for determination of neutron energy spectra, Monte Carlo reconstruction presented by K. Rezac)

EXPERIMENTAL SET-UP AND DIAGNOSTICS

Experimental Set-up



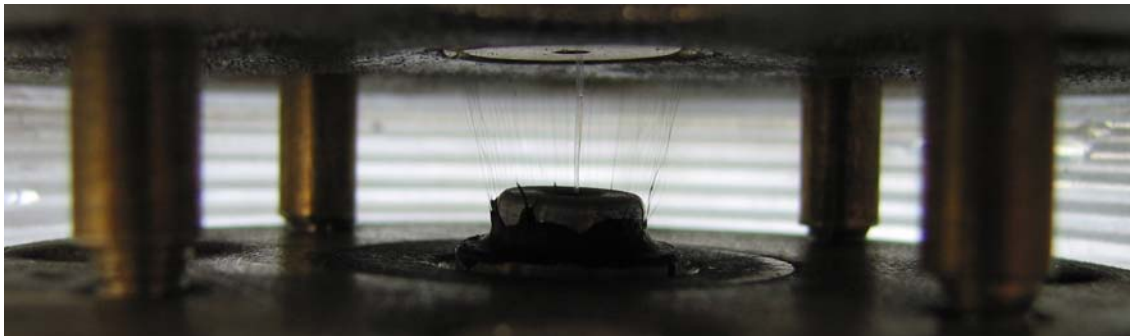
S-300 facility

Marx generator: 8 modules

Electric current: 2.0 MA

Rise-time: 100 ns

Experimental Set-up



Mass percentage of ions in Z-pinch load:
Tungsten 73%, carbon 20%, and deuterium 7%.

S-300 facility

Marx generator: 8 modules

Electric current: 2.0 MA

Rise-time: 100 ns

30 tungsten wires:

$\phi = 7 \mu\text{m}$, $L = 7 \text{ mm}$,

$D_{\text{Anode}} = 10 \text{ mm}$, $D_{\text{C}} = 7 \text{ mm}$

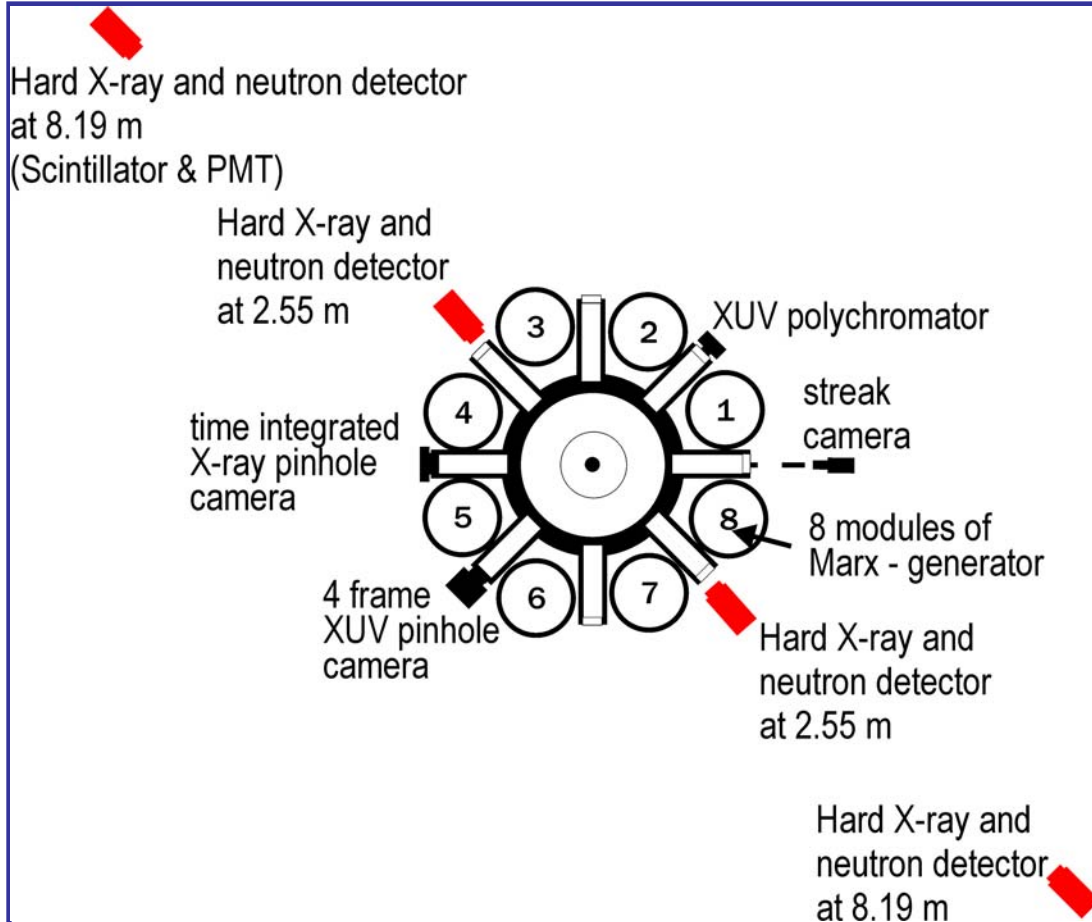
$(\text{CD}_2)_n$ fibre:

$\phi = 100 \mu\text{m}$, $L = 7 \text{ mm}$

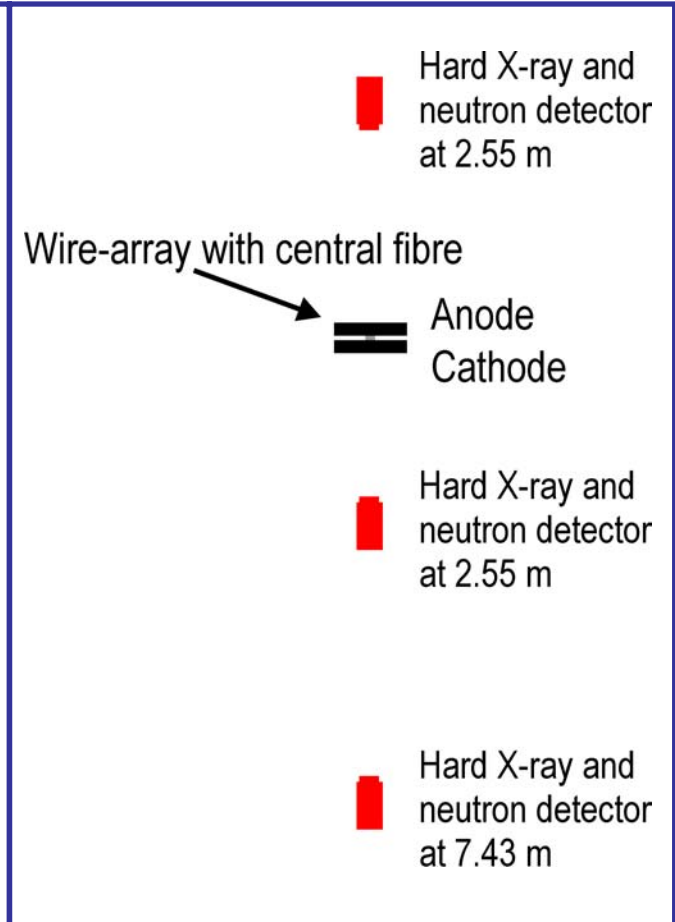
Layout of Diagnostics



End-on view

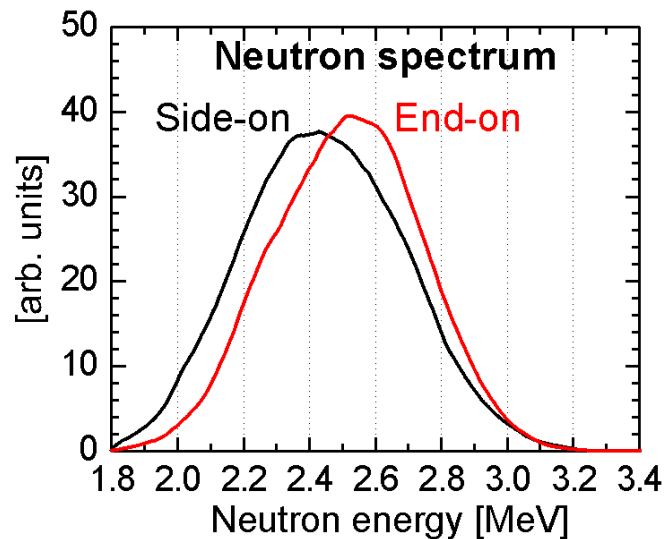
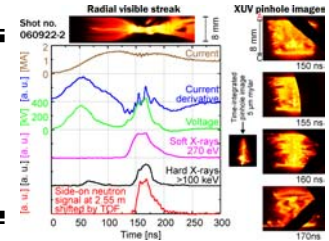


Side-on view



EXPERIMENTAL RESULTS

Neutron Energy Spectra, shot 22-2



Side-on direction (perpendicularly to Z-pinch axis):

mean energy of neutrons about 2.4 MeV,
broader energy spectra radially than downstream

Downstream (axial direction behind cathode):

mean energy of neutrons shifted to 2.55 MeV

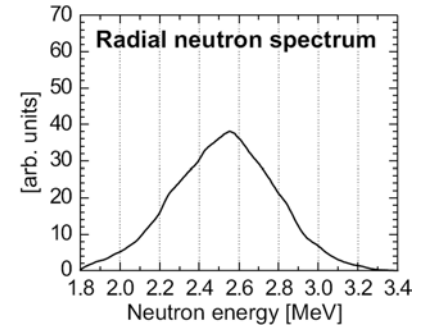
Neutron Energy Spectra, 15 shots



Side-on direction

peak neutron energy: 2.48 ± 0.05 MeV

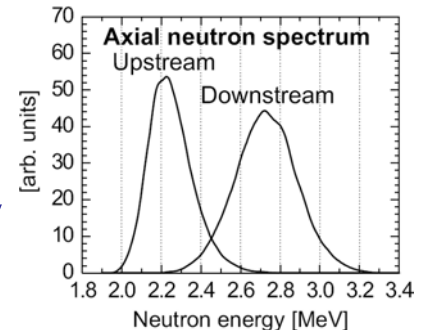
FWHM of energy spectrum: 450 ± 100 keV



End-on direction (downstream)

peak neutron energy: 2.65 ± 0.10 MeV

FWHM of energy spectrum: 350 ± 100 keV

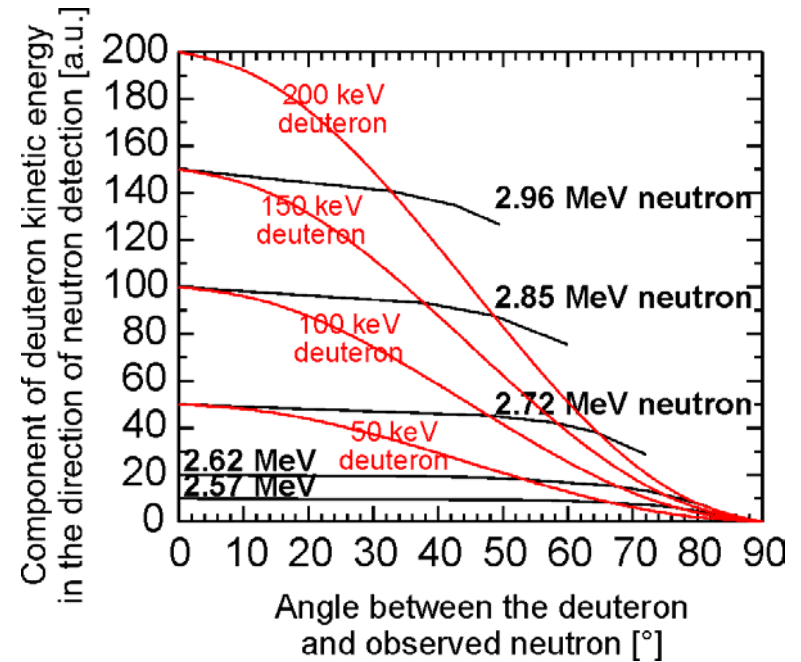
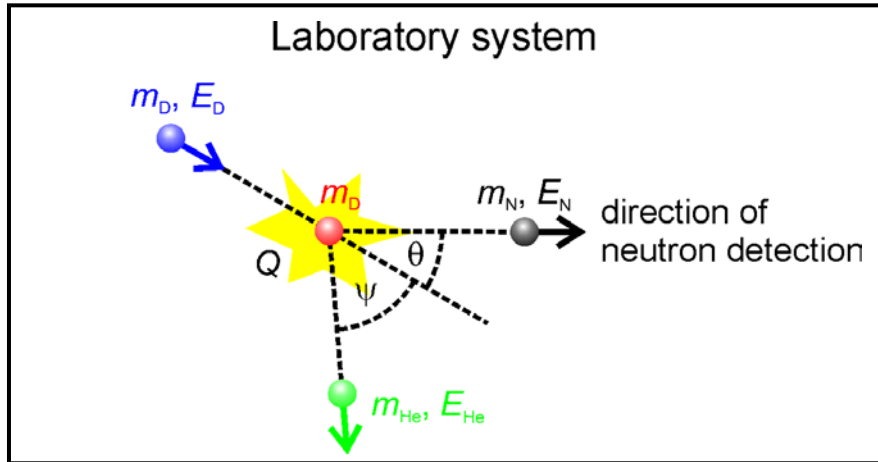


Similar values were observed in other Z-pinch configurations

Compressional Z-pinches: Andrianov 1958, Anderson 1958, Mather 1958...

Plasma foci: Bernstein 1972, Bernard 1975, ...

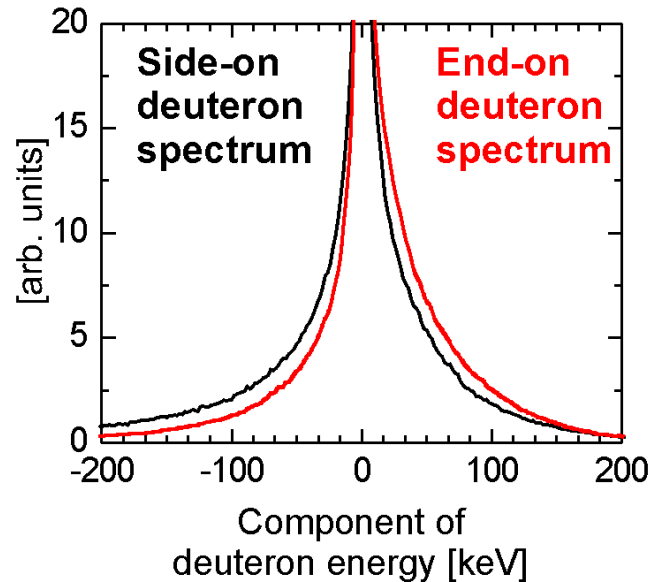
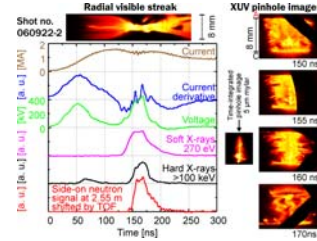
Deuteron Kinetic Energy



$$E_N(E_D, \theta) = E_D \frac{m_N}{2(m_N + m_{He})} \cdot \left(\cos \theta + \sqrt{\frac{m_{He}}{m_N} \left(1 + \frac{2Q}{E_D} \right) - \sin^2 \theta} \right)^2 \approx 2.5 \text{ MeV} + f(E_D \cos^2 \theta)$$

Neutron energy is given mainly by the component of deuteron kinetic energy in the direction of neutron detection.

Energy Distribution Function of Reacting Deuterons, shot 22-2



Axial direction: mean energy component

$$\langle |E_{\parallel}| \rangle = 60 \text{ keV}$$

Side-on direction: mean energy component

$$\langle |E_{\perp}| \rangle = 40 \text{ keV}$$

Fusion neutrons were produced mainly by deuterons

with the average kinetic energy $\langle E_d \rangle = \langle E_{\parallel} \rangle + 2\langle |E_{\perp}| \rangle \doteq 150 \text{ keV}$

Anisotropy of neutron flux: below 1.2

Trajectories of Fast Deuterons



- Neutron emission anisotropy (2.7 MeV neutrons downstream)
implies the role of beam-target model
- Broad neutron energy spectrum in the radial direction
100 keV deuterons in the radial direction
indicate that the linear beam-target model did not occur

- Generalized beam target model

Trajectories of fast deuterons could be complex

E.g. trajectories described by Bernstein 1972

It could explain:

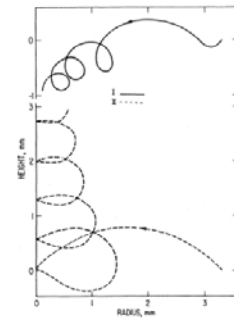
observed neutron spectra

neutron flux anisotropy (<1.5)

Problem: how are these deuterons accelerated?

emission during expansion phase

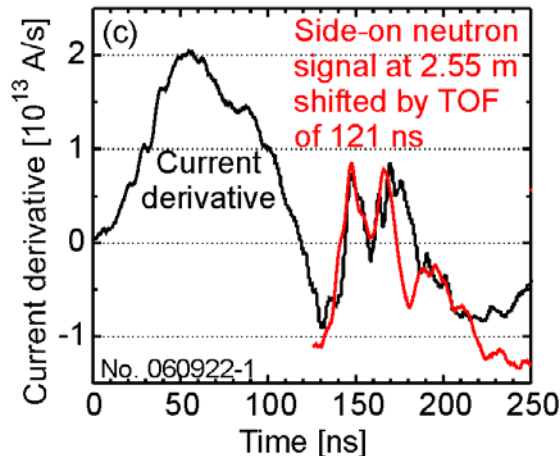
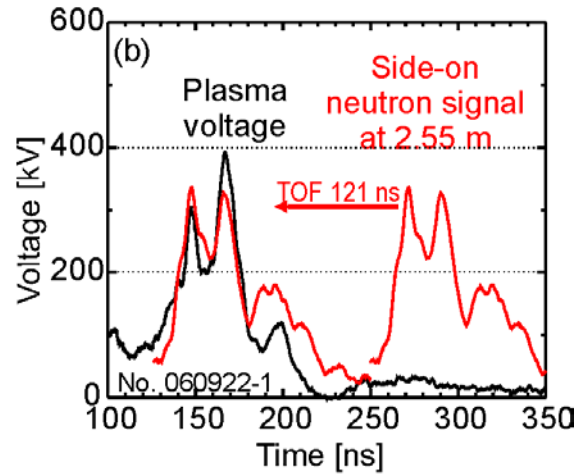
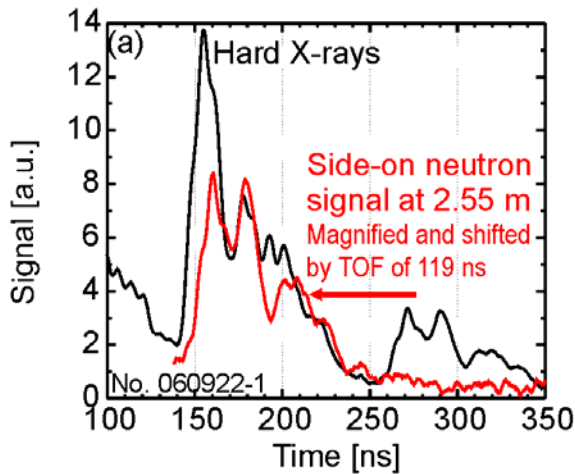
relatively long duration



Temporal correlations



!!! Side-on neutron detector as close to the neutron source as possible !!!

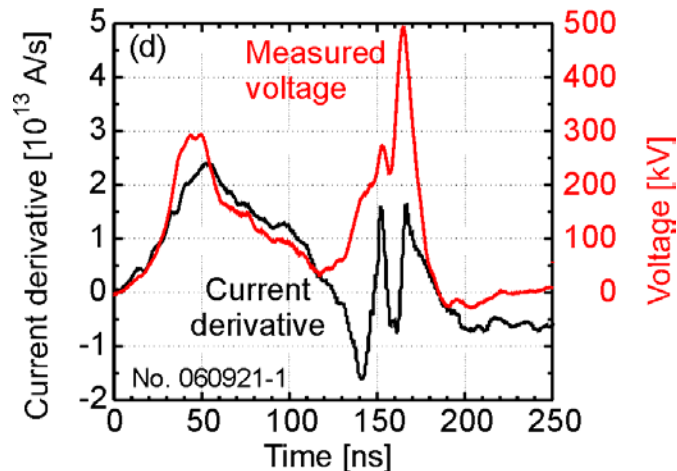


Evident correlation between side-on neutron signal and HXR, voltage, and dI/dt waveforms.

$$V = \frac{d}{dt} \left((L_P + L_0)I \right) + R_P I = (L_P + L_0) \frac{dI}{dt} + \left(\frac{dL_P}{dt} + R_P \right) I$$

Induced voltage LdI/dt contributed to voltage peak 118 ns: time-of-flight of 2.45 MeV neutrons to detector located at 2.55 m

Plasma Resistance and Inductance



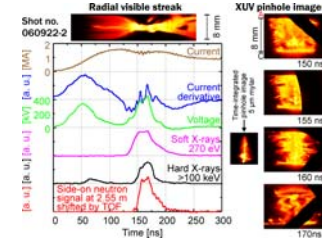
!!! Measurement of voltage V , current derivative dI/dt and current I enabled the estimation of the $R_P + dL_P/dt$ term from

$$R_P + \dot{L}_P = \frac{V - L\dot{I}}{I}$$

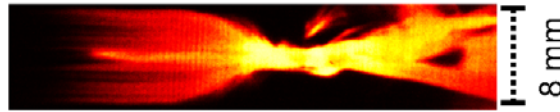
$L = L_P + L_0$ is the inductance including the external inductance of transmission line L_0

We assumed that the inductance $L(t)$ of about 9 nH was approximately constant during the implosion.

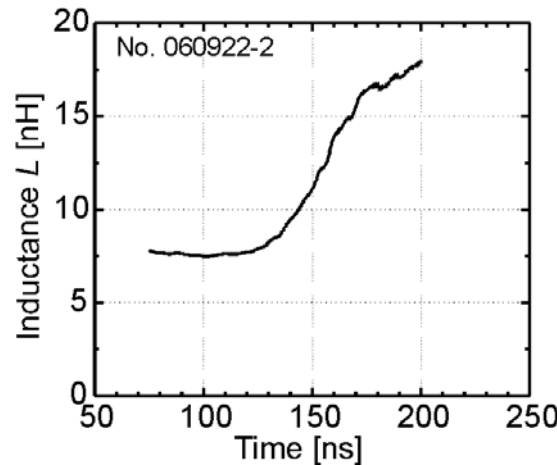
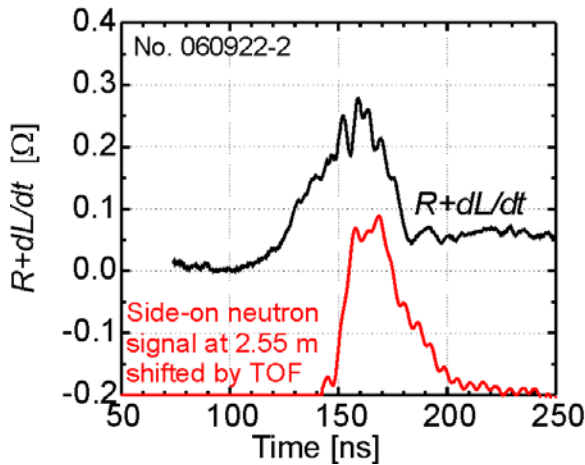
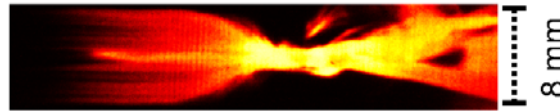
Plasma Resistance and Inductance



Radial visible streak



Radial visible streak



Energy input

$$\int (R_P + \dot{L}_P) I^2 dt = 20 \text{ kJ}$$

During implosion:

0.2 Ω could be ascribed to dL_P/dt

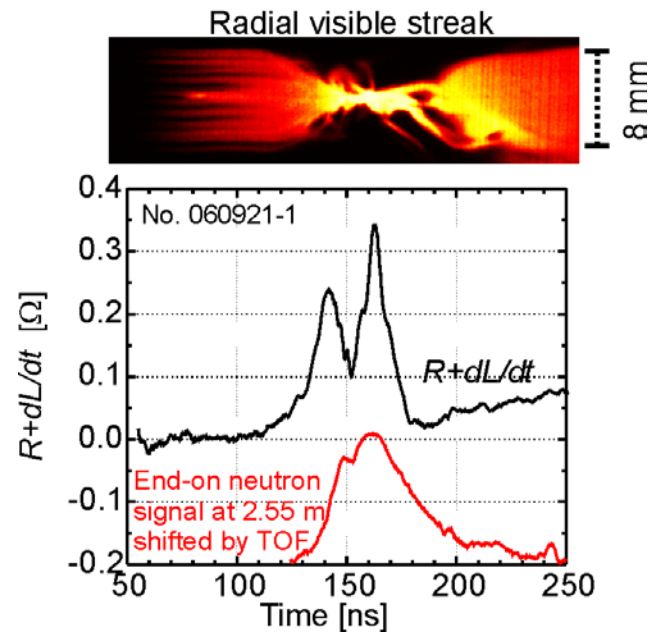
$$|\dot{L}_P| = \frac{\mu l v_{\text{imp}}}{2\pi R}$$

If plasma resistance is neglected:

1. No evidence of the plasma expansion
2. Increase of inductance to 18 nH: **collapse to 10 μm diameter OR increase of effective plasma length**
3. Magnetic energy $\frac{1}{2} \Delta L_P I^2 \doteq 10 \text{ kJ}$ stored after the plasma radiation

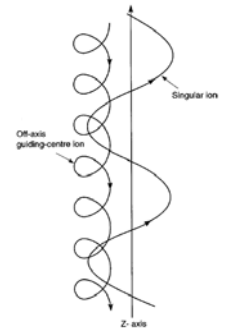
Plasma Resistance and Inductance

In post-stagnation phase: significant role of enhanced resistance
(Spitzer is not sufficient)



The value of $0.2 \div 0.4 \Omega$ also measured by Bernard 1978 and Decker 1983
Energy conservation – debated issue in Z-Pinches (Whitney04, Haines06)

Acceleration of Deuterons



Experimental results best fit to generalized beam-target model.

However, it is not clear how the beam of fast deuterons was created?

- **Diode action:** Because of the correlation between neutron emission and 400 kV voltage peaks, deuterons could be accelerated to 150 keV energies by diode action in peripheral plasma. Axially accelerated deuterons have to be bent by magnetic fields.

Problem: the origin of plasma voltage during post-stagnation phase

- **Microturbulences:** High plasma voltage during the post-stagnation phase could be a result of microturbulences in plasma (Haines 2001). Microturbulences could not only induce high electric fields but also form a high energy tail of the ion distribution function (Ryutov 2000).

Advantages:

1. radial velocities of deuterons, neutron flux anisotropy
2. neutron emission during the stagnation and expansion
3. measured enhancement of plasma resistance

Summary



Implosion of tungsten wire-array onto deuterated fibre:

- Neutron yield of 10^9 during plasma stagnation at fibre
- Reconstruction of neutron energy spectra
- Estimation of kinetic energies of deuterons producing fusion reactions
- Measurement of voltage and plasma resistance and time-varying inductance
- Discussion of acceleration mechanisms

Suggestions for PF-1000



- To have one neutron detector as close to the neutron source as possible in order to measure neutron emission time.
- To measure side-on neutron emission in order to estimate radial velocity distribution of deuterons.
- To measure plasma voltage in order to evaluate the energy input and plasma resistance.

Acknowledgements



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